Detectability of intermediate mass binary black holes



Juan Calderón Bustillo

Center for Relativistic Astrophysics

Georgia Institute of Technology



in collaboration with: P. Laguna, D.Shoemaker, Karan P. Jani, Tito dal Canton, Francesco Salemi & Ian Harry.

7th Iberian Gravitational Waves Meeting

Bilbao, Spain, May 2017

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Introduction: matched filter searches (or modelled searches)



GW151226 (or Boxing Day event)

LSC+ Virgo, PRL 116, 141103 (2016)

Key facts of modelled searches.

- GW signals are normally weaker than the detector background noise.
- Compare (via matched filter) the detector output with waveform models (or templates) —> Signal-to-noise ratio.
- Templates need to be accurate representations of the GW signal.
- Wrong modelling leads to loss in SNR and damages the detection process.



LSC + Virgo, PhysRevX.6.041015

Low mass ratio q=m1/m2 and total mass below M=m1+m2=100Msun

GW150914



LSC + Virgo, PhysRevLett.116.241102

Data is uninformative about precession as the posterior is almost equal to the prior



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LSC + Virgo, Phys. Rev. Lett. 116, 241103

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Search for intermediate mass black hole binaries in the first observing run of Advanced LIGO

During their first observational run, the two Advanced LIGO detectors attained an unprecedented sensitivity, resulting in the first direct detections of gravitational-wave signals produced by stellar-mass binary black hole systems. This paper reports on an all-sky search for gravitational waves (GWs) from merging intermediate mass black hole binaries (IMBHBs). The combined results from two independent search techniques were used in this study: the first employs a matched-filter algorithm that uses a bank of filters covering the GW signal produced by the second is a generic search for GW transients (bursts). No GWs from IMBHBs were detected; herefore, we constrain the rate of several classes of IMB intergrave Element energy obtained for black holes of individual mass 100 M_{\odot} , with spins aligned with the binary orbital angular momentum. For such systems, the merger rate is constrained to be less than 0.93 Gpc⁻³ yr⁻¹ in comoving units at the 90% confidence level, an improvement of nearly two orders of magnitude over previous upper limits.

5,061102 (2016)

LSC + Virgo,arXiv 1704.04628

Motivations summary:

- No detections with total mass above 100Msun.
- All of them showed mass ratios below ~3.5
- No clear signs of precession.

- Black holes with total mass M > 100Msun. \bullet
- If they exist, these would be the strongest CBC-like source for gravitational waves. \bullet



LSC + Virgo,arXiv 1704.04628

Challenges:

- Signal dominated by its merger and ringdown parts.
- Higher modes impact the signal (coming later)
- Templates not as accurate as for the case of lower mass systems.
- Background noise can reproduce easily these signals (Sine gaussian glitches, A.Torres' talk)

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What our (modelled) searches lack: Higher modes & Precession.



Non-Precessing System

Precessing System



C.Capano, LIGO-G1601021, GR21 conference.

- Presence of in-plane spin components causes the orbital plane to precess.
- This causes modulations in the amplitude and frequency of the signal.
- The longer the signal the more precession will hurt (similar to eccentricity).
- Currently used for parameter estimation (SEOBNRv3, PhenomP) but not for searches.



Impact of precession

JCB, Laguna, Shoemaker, PRD, ArXiv1612.02340v2

Introduction: Higher order modes

JCB+, Phys. Rev. D 93, 084019



Harmonics: Depend on orientation

$$h_p(\Xi; r, \theta, \varphi; t) = h_+ - ih_\times = \frac{1}{d_L} \sum_{\ell \ge 2} \sum_{m=-\ell}^{m=\ell} Y_{\ell,m}^{-2}(\theta, \varphi) h_{\ell,m}(\Xi; t)$$

Modes: Depend on source parameters

Impact of higher order modes



JCB, Dal Canton +, In prep. LIGO-G1700062-v6





JCB, Husa, Sintes, Puerrer PRD 93.084019 (2016)

What do we pay?

• Template banks are designed such that not more than **10% of signals are missed** (at a given SNR), this requires a **minimum overlap of 0.965** to targeted signals.



 Precession vs. higher modes: SEOBNRv2 (Taracchini +'12,'14, Pürrer 16) trying to recover Georgia Tech Numerical Relativity waveforms (Jani+ 16).



q=10, non-spinning only higher modes

JCB, Laguna, Shoemaker, Accepted in PRD , ArXiv1612.02340v2

 Precession vs. higher modes: SEOBNRv2 (Taracchini +'12,'14, Pürrer 16) trying to recover Numerical Relativity waveforms (Jani+ 16).



q=10, non-spinning only higher modes q=7, precessing + higher modes

JCB, Laguna, Shoemaker, Accepted in PRD , ArXiv1612.02340v2

 This will cause losses larger than 10% for q > 4 sources when M > 100 Msun for aligned spin cases. (Varma + 14, JCB + 15, JCB +16).



JCB, Husa, Sintes, Puerrer PRD 93.084019 (2016)





JCB, Laguna, Shoemaker, PRD , ArXiv1612.02340v2

Isolating precession: face-on losses.



q > 3 sourcesprecessingDashed: strong

precession (chi_p > 0.5)

GW150914 - like sources

JCB, Laguna, Shoemaker, PRD , ArXiv1612.02340v2

Signals omit higher modes: GOOD separation from background triggers Signals include higher modes: POOR separation from background triggers



JCB, Dal Canton +, In prep. LIGO-G1700062-v6

- The signal-to-noise ratio is the reference value for estimating the significance of triggers.
- However, we must discriminate terrestrial glitches from real GW.
- One of the main ways of doing this is the so called chi-square value
- Mismatch between GW and template, will reduce SNR *and* increase chi-square.

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JCB, Dal Canton +, In prep. LIGO-G1700062-v6

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$$\Phi_{\rm IMR} \sim \begin{cases} \frac{3}{128\eta} \sum_{j=0}^{7} \varphi_j(\Xi) (\pi M f)^{(j-5)/3} + \dots & f < f_1, \\ \frac{1}{\eta} \left(\dots + \beta_2 \log(f) - \frac{1}{3} \beta_3 f^{-3} \right) & f_1 \le f < f_2 \\ \frac{1}{\eta} \left(\dots + \alpha_2 f^{-1} + \frac{4}{3} \alpha_3 f^{3/4} + \alpha_4 \tan^{-1}(af+b) \right) & f \ge f_2 \end{cases}$$





LSC+Virgo, PhysRevLett.116.221101 GW150914 TestingGR paper

$$\Phi_{\rm IMR} \sim \begin{cases} \frac{3}{128\eta} \sum_{j=0}^{7} \varphi_{j}(\Xi) (\pi M f)^{(j-5)/3} + \dots & f < f_{1}, \\ \frac{1}{\eta} \left(\dots + \beta_{2} \log(f) - \frac{1}{3} \beta_{3} f^{-3} \right) & f_{1} \le f < f_{2} \\ \frac{1}{\eta} \left(\dots + \alpha_{2} f^{-1} + \frac{4}{3} \alpha_{3} f^{3/4} + \alpha_{4} \tan^{-1}(af+b) \right) & f \ge f_{2} \end{cases}$$



Test Case: q=8, M=170Msun BBH, edge-on, SNR=50



- Absence of higher modes can lead to <u>apparent fake violations of GR</u>
- The effect of higher modes is accounted by the alpha and beta corrections to GR.
- Larger impact at large post-Newtonian orders
- Wrong signal recovery when GR corrections are switched off.

On its way: implementing a search including higher order modes

Search including higher order modes:

- We build on the existing pyCBC search algorithm.
- Higher modes force us to include two extra parameters: the <u>orientation</u> <u>angles</u> of the source.
- Polarisation can be analytically maximised over.
- Growth of the number of required templates.

Template bank including higher modes: used EOBNRv2HM model (Pan +, PhysRevD.81.084041)



- All <u>LIGO detections</u> so far are characterised by a <u>low mass ratio</u>, total mass <u><</u> <u>100Msun</u>, and <u>no clear evidence for a precessing</u> orbital plane.
- Omission of physical effects like <u>higher modes or precession damages our</u> <u>ability to detect</u> binary black holes, specially those not satisfying the above conditions.
- Precession can have a large impact even for large total masses.
- Losses up to 30% occur for the current version of Advanced LIGO.
- Need for waveform models including the effects of both precession and higher order modes (Ongoing work by Khan, Mishra, Cotesta, UIB group +).

Ongoing work:

- Development of a search using such waveform model, and hopefully discover IMBBH's?.
- We need a waveform model including both precession and higher order modes.

$$\rho = \frac{|\langle h|s\rangle|}{\sqrt{\langle h|h\rangle}} \qquad \langle a|b\rangle \equiv 4 \int_0^\infty \frac{\tilde{a}^*(f)\tilde{b}(f)}{S_n(f)} \mathrm{d}f$$

- SNR is just the projection of the signal s onto the template h.
- If h does not accurately represent the GW contained in s, SNR will be suboptimal.

$$O(a|b) = \frac{\langle a|b \rangle}{\sqrt{\langle a|a \rangle \langle b|b \rangle}} \in [0,1]$$

$$\rho(a|b) \sim O(a|b)$$

$$\rho \sim 1/D_{Hor}$$

$$O(a|b) = [0,1]$$

$$D_{Hor} \sim O(a|b)$$

$$D_{Hor} \sim O(a|b)$$

- Inaccurate templates lead to a loss of SNR and sensitive volume.
- Omission of physical effects causes templates to be inaccurate: event loss and parameter bias

- We look for weak signals in strong noise: usage of <u>matched filter</u> technique.
- We use <u>templates</u> as filters for the incoming signal, and compute its <u>signal-to noise ratio (SNR)</u>.
- <u>Noise is non-Gaussian</u>: we need to discriminate noise transients from real GW's.







Bonus: Parameters of Binary Black Holes' signals



• 2 Masses:

$$M = m_1 + m_2 \qquad q = \frac{m_1}{m_2} \ge 1$$
• 6 Spin Components:

$$\chi_{eff} = \frac{s_1^z m_1 + s_2^z m_2}{m_1 + m_2}$$
2 Spin Components
• Orientation of the Binary (2 angles)
• Orientation of the Detector (3 angles)
• Distance
• Time of arrival
• Develoal Parameters

Analytical

Maximisation

over angles



No Higher

Modes



No in-plane spins

Bonus: Significance of GW150914



LSC + Virgo, Phys. Rev. D 93, 122003

- Need to define some ranking statistic.
- Compute the distribution of the ranking statistic in background: time-slides method.

$$\hat{\rho} = \begin{cases} \rho \left[(1 + (\chi_r^2)^3)/2 \right]^{-\frac{1}{6}}, & \text{if } \chi_r^2 > 1, \\ \rho, & \text{if } \chi_r^2 \le 1. \end{cases}$$

Bonus: Significance of GW150914



LSC + Virgo, Phys. Rev. X 6, 041015 (2016)





Key facts:

- EOBNRv2HM (non-spinning).
- Inclination and azimuth (theta and phi below) are now template parameters (polarisation is maximised over)
- Size ~6000 templates. Modifying now pycbc_inspiral_skymax to use this.







- Lower impact of higher modes due to signals lasting longer in band.
- Rather constant impact of both precession and HM as a function of total mass.
- Slightly larger impact of precession due to waveforms being longer

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