



Beyond the Unitarity Bound in $AdS / CFT_{(A)dS}$

Christoph Uhlemann

Institut für Theoretische Physik und Astrophysik Universität Würzburg

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[JHEP01(2012)123 Tomás Andrade, CFU]

Introduction

AdS/CFT and beyond

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▷ many more brane configurations, gauge/gravity etc.

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Boundary metric (A)dS: [Marolf, Rangamani, van Raamsdonk '11] ▷ curved-space QFT on boundary ▷ branes ending on branes [Aharony, Berdichevsky, Berkooz, Shamir '11] ▷ multi-layered AdS/CFT

$$\int \mathcal{D}\hat{\phi}\Big|_{\hat{\phi}|_{\partial\mathcal{M}} = \delta\phi} e^{iS_{\text{sugra}}[\hat{\phi}]} = \langle e^{i\int_{\partial\mathcal{M}}\delta\phi \,\mathcal{O}} \rangle_{\text{SCFT}}$$

$$\Delta_{\mathcal{O}}(\Delta_{\mathcal{O}} - d) = m_{\phi}^2 =: -d^2/4 + \nu^2 \quad \text{ i.e. } \Delta_{\mathcal{O}} = d/2 \pm \nu$$

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 $\begin{array}{lll} \Delta_{\mathcal{O}} \in \mathbb{R} & \longleftrightarrow & \mathsf{BF} \text{ bound } m_{\phi}^2 \geq -d^2/4 =: m_{\mathrm{BF}}^2 & \checkmark \\ \text{unitarity bound } \Delta_{\mathcal{O}} \geq d/2 - 1 & \longleftrightarrow & \mathsf{Neumann \ b.c. \ only \ for \ } \nu^2 \leq 1 & \checkmark \end{array}$

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 \triangleright ghosts on global AdS

[Andrade, Marolf '11]

▷ 2-point function sick on Poincaré AdS

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Here: set up holographic renormalization for CFTs on $\rm (A)dS$ study unitarity holographically

Outline:

 \triangleright CFT on AdS from the holographic perspective:

- the geometry
- regularization and renormalization
- beyond the unitarity bound
- multi-layered AdS/CFT
- \triangleright CFT on dS
 - geometry
 - beyond the unitarity bound

$$ds^{2} = -(1+\rho^{2})dt^{2} + \frac{d\rho^{2}}{1+\rho^{2}} + \rho^{2}d\Omega_{d-1}^{2}$$



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$$d\zeta^{2} + \sin^{2}\zeta d\Omega_{d-2}^{2}$$





$$\begin{split} (\rho,\zeta) \to (R,z) &: \quad \rho^2 = \csc^2 z \, \cosh^2 \! R - 1 \\ \rho^2 {\sin^2} \zeta = \cot^2 z \, \cosh^2 \! R \end{split}$$

AdS_{d+1} in global coordinates:

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$$\rightarrow ds^2 = dR^2 + \cosh^2 R \, ds^2_{\mathsf{AdS}_d}$$
$$ds^2_{\mathsf{AdS}_d} = \frac{1}{\sin^2 z} \left(dz^2 - dt^2 + \cos^2 z \, d\Omega^2_{d-2} \right)$$

conformal boundary at $|R| = \infty$ is $2 \times \text{AdS}_d$



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conformal boundary at $|R| = \infty$ is $2 \times \text{AdS}_d$ $\rightarrow \mathbb{Z}_2$ orbifold for single copy







Divergences and Renormalization

Klein-Gordon field
$$\phi$$
 on AdS_{d+1}/ \mathbb{Z}_2 , $y := e^{-R}$, $m_{\phi}^2 = -d^2/4 + \nu^2$:



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$$\begin{split} \chi : \Box_{\mathsf{AdS}_d} \chi &= M^2 \chi \qquad M^2 = -(d{-}1)^2/4 + \mu^2 \\ \rightarrow \phi &= \chi(z,t,\Omega_{d-2})\,f(y) \end{split}$$



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$$\delta S = \mathsf{EOM} + \int_{y=1}^{z} \delta \phi \partial_y \phi + \int_{y=0}^{z} \delta \phi \partial_y \phi + \int_{z=0}^{z} \delta \phi \partial_z \phi$$



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 $\sum_{z=0}^{} \underbrace{y \to 0}_{\operatorname{Dirichlet}_{d+1}/2}$

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 $\delta S \supset \text{divergences from } y/z \rightarrow 0: \text{ regularize } y \ge \epsilon_1, \ z \ge \epsilon_2$ $\rightarrow \partial \mathcal{M}: \ \mathbb{Z}_2\text{-fixed surface } \partial_0 \mathcal{M}, \ y = \epsilon_1/z = \epsilon_2 \text{ surfaces } \partial_{1/2} \mathcal{M}$

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$\mathsf{Dirichlet}_d$ and the unitarity bound

$Dirichlet_d$ beyond the unitarity bound

Choose $\nu \in (1,2)$ and fixed Dirichlet_d . For finite δS :

$$S_{\mathsf{ct}} = -\frac{1}{2} \int_{\partial_1 \mathcal{M}} \left[(d/2 - \nu)\phi^2 + \frac{1}{2(\nu - 1)}\phi \,\Box^W \phi \right] + \frac{1}{4(\nu - 1)} \int_{\partial \partial \mathcal{M}} \phi \mathcal{L}_n \phi$$

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associated inner product: $\langle \phi_1, \phi_2 \rangle = \langle \chi_1, \chi_2 \rangle \cdot \langle f_1, f_2 \rangle_{\mathrm{ren}}$

$$\langle f_1, f_2 \rangle_{\mathrm{ren}} = \langle f_1, f_2 \rangle_{\mathrm{bare}} - \frac{1}{2(\nu-1)} (\cosh R)^{d-2} f_1^* f_2 \big|_{R \to \infty}$$

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Spectrum:
$$\mu = 2n + \frac{1}{2} - a_{D/N}\nu + b_{even/odd}$$
, $||f||^2 = \frac{(n+b/2)!}{\mu\Gamma(2\mu-2n-b)}$
 $\checkmark ||f||^2 > 0$ for Dirichlet_{d+1}
 \checkmark ghosts for $\nu > 1$ Neumann_{d+1}
 \triangleright results extend to generic ν
 \triangleright log-modes for $2\nu \in \mathbb{Z}$

Saturating the unitarity bound: $\nu = 1$

 $\rightarrow \log$ terms in asymptotic expansion

 \rightarrow new finite combination of boundary terms $\propto \kappa$

$$S_{\rm ct} = -\frac{1}{2} \int_{\partial_1 \mathcal{M}} \left[(d/2 - 1)\phi^2 - \left(\log y + \kappa\right)\phi \,\Box^W \phi \right] - \frac{1}{2} \int_{\partial \partial \mathcal{M}} (\log y + \kappa)\phi \mathcal{L}_n \phi$$

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Dirichlet_{d+1}: spectrum straightforward, norms positive Neumann_{d+1}: spectrum only implicitly (dilogs)

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shown for lowest- μ^2 state:

 $\rhd~\mu^2 < 1/2$ generally, BF bound violated for certain κ

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in summary: satisfactory renormalization procedure for $Dirichlet_d$ CFT unitarity violation reflected in bulk by ghost states

Neumann_d: dimensional reduction

$$\begin{split} \phi &= \sum_n \alpha_n \phi_n \text{ , } \quad \phi_n = \chi_{\mu_n} \cdot f_n, \ n \in \mathbb{N} \quad \sim \text{`tower' of } \mathsf{AdS}_d \text{ modes} \\ & \langle \phi_n, \phi_m \rangle = \underbrace{\langle \chi_{\mu_n}, \chi_{\mu_m} \rangle}_{\text{finite for } \text{Dirichlet}_d} \underbrace{\langle f_n, f_m \rangle}_{\rightarrow \text{ finite } \langle f_n, f_m \rangle_{\text{ren}}}_{\text{w/ c.t. on } \partial_1 \mathcal{M}} \end{split}$$

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Neumann_d: χ_{μ_n} normalizable only for $\mu_n^2 < 1$ ($z \rightarrow 0$ divergences)

add c.t. on $\partial_2 \mathcal{M}$ s.t. all $\langle \chi_{\mu_n}, \chi_{\mu_m} \rangle$ finite?



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$$S_{\mathsf{ct}} = \int\limits_{\partial_2 \mathcal{M}} \alpha_1 \phi^2 + \alpha_2 \phi \Box \phi + \alpha_3 \mathcal{R} \phi^2 + \dots$$

need $\alpha_i = \alpha_i(\mu)$, e.g. $\alpha_2 = 1/(1 - \mu_n)$ $\rightarrow w/$ fixed α_i only one $\mu_n^2 > 1$ normalizable



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reduction of (infinite) tower of AdS_d modes to finitely many \sim dimensional reduction, no localization in radial direction

 $\begin{array}{ll} \mbox{Including } \mu^2 \geq 1 \mbox{ modes: } \langle \chi_\mu, \chi_\mu \rangle \mbox{ indefinite } \rightarrow \mbox{ ghosts } & \mbox{[Andrade, Marolf '11]} \\ \rightarrow \mbox{ add no c.t. on } \partial_2 \mathcal{M} \mbox{ s.t. only } \mu^2 < 1 \mbox{ normalizable } \end{array}$

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 $\label{eq:solid} \begin{array}{l} {\rm dashed/solid} \sim {\rm even/odd} \\ ||f||^2 < 0 \mbox{ for } \nu \in (1,2) \ensuremath{\checkmark} \\ ||f||^2 > 0 \mbox{ for } \nu < 1 \ensuremath{\checkmark} \\ \nu & \mbox{ and } \nu \in (2,3)? \end{array}$

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 \triangleright ghost-free Neumann theory for $\nu > 1$

ho truncation breaks conformal invariance ightarrow no unitarity bound \checkmark

$\mathsf{AdS}/\mathsf{CFT} \text{ iterated}$

Prospects for multi-layered AdS/CFT?

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Prospects for multi-layered AdS/CFT?

(Parental advisory: speculative!)



→ 2-fold AdS/CFT: gravity_{d+1} \cong CFT_{d-1} → Bengt Nilsson's talk even generalize to *n*-fold? from scalar: expect dim. reduction of gravity_{d+1} w/ Neumann_d → multiple iterations trivialized

unitarity in $\mathsf{CFT}_{\mathrm{dS}}$ holographically

 AdS_{d+1} in global coordinates:

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 $(\rho, t) \rightarrow (R, \tau)$: $\rho = \cosh(H\tau) \sinh R$ $\tan(t) = \sinh(H\tau) \tanh R$

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- \vartriangleright renormalization straightforward, cut-off on R sufficient
- $\triangleright \ \nu \in (1,2)$: spectrum and norms straightforward, $\nu = 1$ more work

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- \triangleright CFT unitarity violation reflected by ghosts immediately \rightarrow Poincaré subtleties can not be blamed on horizon

Conclusion

Unitarity CFT_{AdS}

- \triangleright Dirichlet_d: boundary non-unitarity reflected in bulk by ghost states
- $\,\triangleright\,$ Neumann_d: $\,\cdot\,$ reduced spectrum of normalizable solutions, ghost-free for certain $\nu>1$
 - conformal invariance broken \rightarrow unitarity bound does not apply

 $\,\triangleright\,$ nested AdS/CFT: trivialized by dimensional reduction in this setting

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Thanks for your attention!