# Supergraviti**ES** in d = 11Nucl. Phys. B **855** (2012) 308-319

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3 Subtle links between Supergravity and Chern–Simons Theories in d = 11



## Contents



- 2 An Eccentric but Sensible Theory: Chern–Simons Supergravity
- 3 Subtle links between Supergravity and Chern–Simons Theories in d = 11

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4 Further Work & Questions

The Uniqueness of Supergravity in d = 11

• SUGRA in d = 11 is unique!!

S. Deser, The Uniqueness of D = 11 Supergravity, hep-th/9712064

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• But there are exceptions! It is necessary "to read the fine print"!!

The Uniqueness of Supergravity in d = 11

• Underlying the uniqueness, there are some hypotheses which are very natural but not unbreakable.

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• A simple way of breaking these hypotheses is through a **Chern–Simons Theory**.

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## In Further Work & Questions

## Chern–Simons SUGRA

Chern–Simons theories in different odd dimensions started in the 80's-90's with the work of

A. Achúcarro,

M. Bañados,

A. Borowiec,

A. H. Chamseddine,

M. Ferraris,

M. Francaviglia

M. Henneaux

P. K. Townsend,

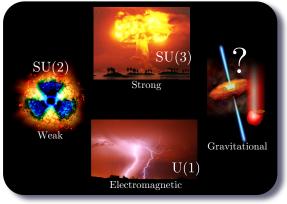
R. Troncoso,

J. Zanelli,

etc...

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## Gauge Theories X-treme



Gauge theories do work.

What about a Supergravity theory where *every* field, *all* of them, belong only to a 1-form *connection*?

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## Gauge Theories X-treme

$$\boldsymbol{A} = \frac{1}{2}\omega^{ab}{}_{\mu}\mathrm{d}x^{\mu}\boldsymbol{J}_{ab} + \frac{1}{l}\boldsymbol{e}^{a}{}_{\mu}\mathrm{d}x^{\mu}\boldsymbol{P}_{a} + \frac{1}{\sqrt{l}}\bar{\psi}_{\mu}\mathrm{d}x^{\mu}\boldsymbol{Q} + \mathrm{extra bosonic terms.}$$

A really daring idea!

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## Gauge Theories X-treme

And the Lagrangian??

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## Gauge Theories X-treme

$$\langle F \wedge *F \rangle$$

F. Izaurieta - arXiv: 1103.2182 (hep-th) - Fizaurie@ucsc.cl SUGRAs in d = 11 - February 2, 2012

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## Gauge Theories X-treme

 $\langle F \wedge *F \rangle \oslash$ Hodge dual  $* \Rightarrow$  a metric field  $g_{\mu\nu}(x) \notin A!!$ 

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## Gauge Theories X-treme

### :-( And then?

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## Gauge Theories X-treme

#### Chern–Simons Lagrangian Form

$$\mathcal{L}_{\mathrm{CS}}^{(2n+1)}(\boldsymbol{A}) = \kappa(\boldsymbol{n}+1) \int_{\tau=0}^{\tau=1} \mathrm{d}\tau \left\langle \boldsymbol{A} \wedge \left(\tau \mathrm{d}\boldsymbol{A} + \tau^2 \boldsymbol{A} \wedge \boldsymbol{A}\right)^{\wedge \boldsymbol{n}} \right\rangle$$

 $\langle \mathbf{T}_{A_1} \cdots \mathbf{T}_{A_{n+1}} \rangle = (\text{Super}) \text{Algebra Invariant Tensor}$ (Close cousin of the superalgebra killing metric)

# Gauge Theories X-treme

#### Advantages of CS Theories

• The Lagrangian has a deep mathematical meaning (Chern–Weil)!

- **2** Genuine **gauge** theories, in fiber bundle sense.
- Off-shell invariant (without auxiliary fields).
- They *seem* to be renormalizable (only proven in d = 3).
- In general, no boson-fermion matching.
- It does include gravity (Lanczos–Lovelock)!
- O No ghosts.
- No higher spin fields.
- **2** Propagating degrees of freedom in d > 3.
- Background-free.

# Gauge Theories X-treme

#### Disadvantages of CS Theories

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- It works only in **odd** dimensions!
- 2 Each individual term is not explicitly invariant.
- So The number of terms grows exponentially with *d*.
- **④ Highly non-linear** in d > 3!
- S Irregular constraint system (complex dynamics).
- Quantum behavior is **unknown** in d > 3.

## Gauge Theories X-treme

An odd number...?

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## Gauge Theories X-treme



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# 3 Subtle links between Supergravity and Chern–Simons Theories in d = 11

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## $\mathfrak{osp}(32|1)$ -Chern–Simons SUGRA in d = 11

Chern–Simons Lagrangian in d = 11

$$\mathcal{L}_{\text{CS}}^{(11)}\left(\boldsymbol{A}\right) = 6\kappa \int_{\tau=0}^{\tau=1} \mathrm{d}\tau \left\langle \boldsymbol{A} \wedge \left(\tau \mathrm{d}\boldsymbol{A} + \tau^{2}\boldsymbol{A} \wedge \boldsymbol{A}\right)^{\wedge 5} \right\rangle$$

 $\langle T_{A_1} \cdots T_{A_6} \rangle$  = Superalgebra Invariant Tensor

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## $\mathfrak{osp}(32|1)$ -Chern–Simons SUGRA in d = 11

$$\boldsymbol{A} = \frac{1}{2}\omega^{ab}\boldsymbol{J}_{ab} + \frac{1}{l}e^{a}\boldsymbol{P}_{a} + \frac{1}{5!l}b^{a_{1}\cdots a_{5}}\boldsymbol{Z}_{a_{1}\cdots a_{5}} + \frac{1}{\sqrt{l}}\bar{\psi}\boldsymbol{Q}$$
  
osp (32|1)-connection 1-form.

R. Troncoso, J. Zanelli, Phys. Rev. D 58 (1998) 101703, hep-th/9710180

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## SUGRAs in d = 11

#### But d = 11 is already inhabited!

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## SUGRAs in d = 11

Standard SUGRA, E. Cremmer, B. Julia and J. Scherk (1978),

$$\begin{aligned} \mathcal{L}_{\text{CJS}}^{(11)} &= -\frac{1}{9!} \frac{1}{4\kappa^2} \epsilon_{a_1 \cdots a_{11}} R^{a_1 a_2} \wedge e^{a_3} \wedge \cdots \wedge e^{a_{11}} + \frac{i}{2} \bar{\psi} \wedge \Gamma_{(8)} \wedge \mathrm{D}\psi + \\ &+ \frac{i}{8} \left( T^a - \frac{i\kappa^2}{4} \bar{\psi} \wedge \Gamma^a \psi \right) \wedge e_a \wedge \bar{\psi} \wedge \Gamma_{(6)} \wedge \psi - \frac{1}{2} F \wedge * F + \\ &+ (*F + b) \wedge (\mathrm{d}A_3 - a) + \frac{1}{2} a \wedge b - \frac{1}{3} A_3 \wedge \mathrm{d}A_3 \wedge \mathrm{d}A_3, \end{aligned}$$

the "only" theory of SUGRA in d = 11 (under several hypotheses).

B. Julia, S. Silva, JHEP 0001 (2000) 026, hep-th/9911035

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## SUGRAs in d = 11

$$\Gamma_{(n)} = \frac{1}{n!} \Gamma_{a_1 \cdots a_n} e^{a_1} \wedge \cdots \wedge e^{a_n}$$
$$a = \frac{i\kappa}{4} \bar{\psi} \wedge \Gamma_{(2)} \wedge \psi,$$
$$b = \frac{i\kappa}{4} \bar{\psi} \wedge \Gamma_{(5)} \wedge \psi.$$

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## SUGRAs in d = 11



#### What to do with two SUGRA theories in d = 11?

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## SUGRAs in d = 11

#### Could there be a link between both SUGRA theories?

Surprisingly hard to answer!

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## SUGRAs in d = 11

## In particular, could it be possible to consider standard SUGRA as a sector of a Chern–Simons system? (A priori, **it seems to be no reason to think so!**)

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## WARNING!!!

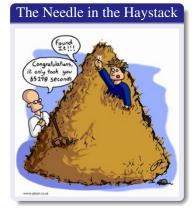
- CS SUGRA is very complicated! Lagrangian with more than a 1000, non-explicitly invariant terms!
- Jumping to compare both of them would be hazardous for mental health!



## WARNING!!!

- The way of finding the "needle in the haystack" is (carefully) using
  - Cartan's Homotopy Formula[1],
  - arguments of dimensional consistency and
  - numerical algoritms[2].

```
[1] hep-th/0603061
[2] 1106.1648 [math-ph]
```



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## CS SUGRA and powers of l

$$\mathcal{L}_{\mathrm{CS}}^{(11)} = \sum_{n=0}^{11} l^{-n} \mathcal{L}_n$$

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Just for the sake of simplicity, a very restricted sector will be explored:

- Order  $l^{-9}$ ,
- Only  $\mathcal{L}_{\text{CS}}^{(11)}\Big|_{b^{a_1\cdots a_5}=0}$  and  $\mathcal{L}_{\text{CJS}}^{(11)}\Big|_{A_3=0}$  will be compared!
- Only the most simple invariant tensor will be used!

## SUGRAs in d = 11

Restricted sector of osp (32|1)-CS Lagrangian:

$$\begin{split} \mathcal{L}_{\text{CS}}^{(11)}\Big|_{\substack{l^{-9}\\b=0}} &= \frac{\lambda}{l^9} \left[ -\frac{1}{4 \times 9! \kappa_2^2} \epsilon_{a_1 \cdots a_{11}} R^{a_1 a_2} \wedge e^{a_3} \wedge \cdots \wedge e^{a_{11}} + \right. \\ &\left. -\frac{7i}{4} \bar{\psi} \wedge \Gamma_{(8)} \wedge \mathcal{D}_\omega \psi - \frac{1}{2} a \wedge b + \right. \\ &\left. + \frac{i}{8} \left( T^a - \frac{i\kappa_1^2}{4} \bar{\psi} \wedge \Gamma^a \psi \right) \wedge e_a \wedge \bar{\psi} \wedge \Gamma_{(6)} \wedge \psi + \right. \\ &\left. + \frac{7}{2^4 \times 5} a \wedge * a - \frac{1}{2^{10} \times 3^2} \frac{(4!)^5}{6!} b \wedge * b \right]. \end{split}$$

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## SUGRAs in d = 11

Standard SUGRA, E. Cremmer, B. Julia and J. Scherk (1978),

$$\begin{aligned} \mathcal{L}_{\text{CJS}}^{(11)} &= -\frac{1}{9!} \frac{1}{4\kappa^2} \epsilon_{a_1 \cdots a_{11}} R^{a_1 a_2} \wedge e^{a_3} \wedge \cdots \wedge e^{a_{11}} + \frac{i}{2} \bar{\psi} \wedge \Gamma_{(8)} \wedge \mathrm{D}\psi + \\ &+ \frac{i}{8} \left( T^a - \frac{i\kappa^2}{4} \bar{\psi} \wedge \Gamma^a \psi \right) \wedge e_a \wedge \bar{\psi} \wedge \Gamma_{(6)} \wedge \psi - \frac{1}{2} F \wedge * F + \\ &+ (*F + b) \wedge (\mathrm{d}A_3 - a) + \frac{1}{2} a \wedge b - \frac{1}{3} A_3 \wedge \mathrm{d}A_3 \wedge \mathrm{d}A_3, \end{aligned}$$

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## SUGRAs in d = 11

 $A_3 = ?$ 

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## SUGRAs in d = 11

$$A_3 \sim \frac{1}{l^3}$$

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## SUGRAs in d = 11

$$A_3 = C_{a_1 \cdots a_5 bc} b^{a_1 \cdots a_5} \wedge e^b \wedge e^c + C_{a_1 \cdots a_5 b_1 \cdots b_5 c} b^{a_1 \cdots a_5} \wedge b^{b_1 \cdots b_5} \wedge e^c + + C_{a_1 \cdots a_5 b_1 \cdots b_5 c_1 \cdots c_5} b^{a_1 \cdots a_5} \wedge b^{b_1 \cdots b_5} \wedge b^{c_1 \cdots c_5}$$

• A<sub>3</sub> seems to be related to the Chern–Simons 3-form for osp(32|1)!!

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## SUGRAs in d = 11

### A deformed version of standard CJS SUGRA in d = 11seems to be contained as a sector of the Chern–Simons Theory!

A priori, no reason to expect such "coincidence"!

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## Further Work & Questions

• It is very likely that standard CJS SUGRA can be considered as an small sector of a Chern–Simons system, when a more general invariant tensor is considered.

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## Further Work & Questions

• What could this relationship between standard SUGRA and Chern–Simons theory in *d* = 11 possibly mean...?

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## Further Work & Questions

 Could a Chern–Simons theory in d = 11 play a rôle in the net of dualities surrounding Strings / M Theory? [1, 2]

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- If it does, to which kind of String theory it could be related?
  - [1] P. Horava, Phys. Rev. D 59 (1999) 046004, hep-th/9712130
  - [2] H. Nastase, hep-th/0306269

## Thank You!



To make this work, we used a bundle of great software, which was *libre* and *gratis*. This software was created by the open source community of mathematicians, physicists, programmers and artists from all around the world. **Thank You!** 

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## Thank You!

